

Towards Probing Ultralight Dark Matter Couplings with Acetylene Spectroscopy

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1) From UDM couplings to SM to frequency variations of AMO resonances

2) Potential from infrared precision measurements with $^{12}\text{C}_2\text{H}_2$
-Modelisation and measurement of a $^{12}\text{C}_2\text{H}_2$ transition at $1.55\ \mu\text{m}$
-Constraints on UDM couplings to SM

3) Potential from frequency measurements of a MW acetylene clock
-Acetylene MW transitions with enhanced sensitivity to μ -variation
-Molecular theory and metrological performances of a MW clock
-Constraints from frequency measurements of a MW acetylene clock

4) Conclusion

1) Variation of fundamental constants from precision measurements

• Why Acetylene?

• Objectives

- State-of-the-art molecular theory and precision measurements

Global Hamiltonian models : Herman and Perry, PCCP 15, 9970 (2013); Lyulin and Perevalov, JQSRT 177, 59 (2016)

Ab-initio theory : Chubb *et al*, JQSRT 204, 42 (2018)

Instrumentations for high-resolution spectroscopy; spectroscopic data; molecular databases HITRAN, GEISA, ...

Cold molecules research : Aiello *et al*, Nat. Commun. 13, 7016 (2022)

- Search for variability of fundamental constants; understand nature of dark matter

Uzan, Living Rev. Relativity 14, 2 (2011)

Safronova *et al*, Rev. Mod. Phys. 90, 025008 (2018)

• Applications of the acetylene molecular theory and spectroscopy

- Atmospheric science

Whitby and Altwicker, Atmos. Environ. 12, 1289 (1978)

- Astrophysics

Didriche and Herman, Chem. Phys. Lett. 496, 1 (2010)

- Frequency metrology

Recommandation Comité Consultatif des Longueurs 1, 2009; Riehle *et al*, Metrologia 55, 188 (2018)

- Probe variations of fundamental constants

Constantin, Vibrational Spectroscopy 85, 228 (2016)

- Photonic-molecular integration

Tharpa *et al*, Opt. Lett. 31, 2489 (2006); Takiguchi *et al*, Opt. Lett. 36, 1254 (2011)

Zektzer *et al*, Laser Photonics Rev. 14, 1900414 (2020)

Fundamental Constants and Ultralight Dark Matter

- The Standard Model and the General Relativity : FC are free parameters of the theory

Uzan, C. R. Physique 16, 576 (2015)

- CODATA 2018 recommended values of fundamental constants

<https://physics.nist.gov/cuu/Constants/index.html>

- Variability from couplings to cosmology and to local fields

- UDM: sub-eV scalar field $\phi(t) = \phi_0 \cos(\omega_\phi t)$; pulsation $\omega_\phi \cong \frac{m_\phi c^2}{\hbar}$; amplitude $\phi_0 = \sqrt{\frac{4\pi G \rho_{DM}}{\omega_\phi^2 c^2}}$

- Galactic halo model : $\rho_{DM} = 0.4 \text{ GeV/cm}^{-3}$; $v \cong 230 \frac{\text{km}}{\text{s}}$; $\sigma_v \cong 10^{-3} c$; $\tau_c = 10^{-4} / \omega_\phi$

Arvanitaki *et al*, PRD 91, 015015 (2015); Stadnik *et al*, PRL 115, 201301 (2015); Freese *et al*, RMP 85,1561 (2013)

=> DM-field induced variation of fundamental constants

- fermion $X=(e,u,d,s)$ masses : $m_X(\phi) = m_X \left(1 + d_{m_X} \cdot \phi / \sqrt{c\hbar/8\pi G} \right)$

- fine structure constant : $\alpha(\phi) = \alpha \left(1 + d_\alpha \cdot \phi / \sqrt{c\hbar/8\pi G} \right)$

- QCD scale parameter : $\Lambda_{QCD}(\phi) = \Lambda_{QCD} \left(1 + d_g \cdot \phi / \sqrt{c\hbar/8\pi G} \right)$

- $\mu(m_e, \Lambda_{QCD}, m_X) = m_p/m_e$: $\mu(\phi) = \mu \left(1 - (d_{m_e} + d_g + 0.036 d_{\hat{m}}) \cdot \phi / \sqrt{c\hbar/8\pi G} \right)$

Damour and Donoghue, PRD 82, 084033 (2010)

=> Sensitivity of atomic, molecular and optical cavity resonance frequency:

$$\frac{\Delta f_{C,A,M}}{f_{C,A,M}} = Q_{\alpha}^{C,A,M} \frac{\Delta \alpha}{\alpha} + Q_{\mu}^{C,A,M} \frac{\Delta \mu}{\mu} + Q_q^{C,A,M} \frac{\Delta(\hat{m}/\Lambda_{QCD})}{\hat{m}/\Lambda_{QCD}}$$

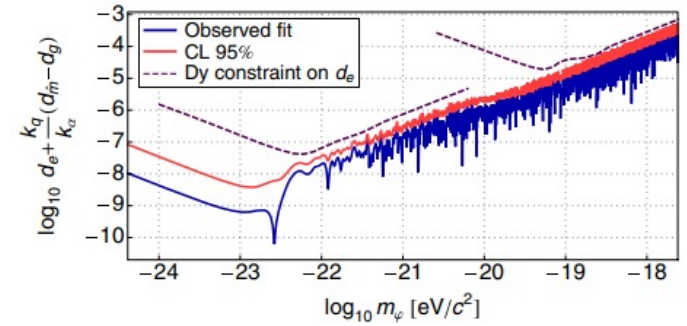
- atomic optical transition $f_{A,opt} = C_{A,opt} \frac{\alpha^2 m_e c^2}{4\pi\hbar} \cdot F_{opt}(\alpha)$
- atomic hyperfine transition $f_{A,HF} = C_{A,HF} \frac{\alpha^2 m_e c^2}{4\pi\hbar} (g_i(m_q/\Lambda_{QCD}) \times m_e/m_p) \cdot F_{HF}(\alpha)$
- vibrational molecular transition : $f_{vib} = C_{vib} \frac{\alpha^2 m_e c^2}{4\pi\hbar} \sqrt{m_e/m_p}$
- rotational molecular transition : $f_{rot} = C_{rot} \frac{\alpha^2 m_e c^2}{4\pi\hbar} m_e/m_p$
- optical cavity resonance : $f_C = C_C \frac{\alpha m_e c^2}{\hbar} \cdot F_C(\alpha, m_X, \Lambda_{QCD})$

• Searches for UDM with atoms and molecules

=> Atomic clocks comparisons

- oscillations from scalar fields of Rb/Cs clocks

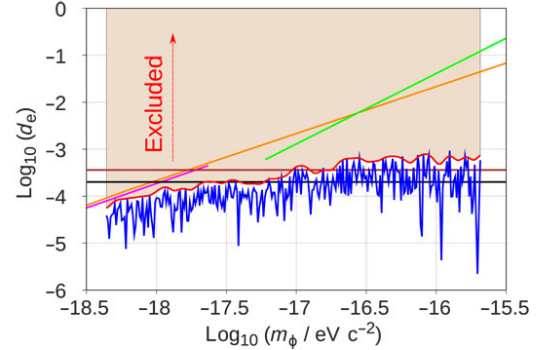
Hees *et al*, Phys. Rev. Lett. 117, 061301 (2016)



=> Atomic clocks/cavity comparisons

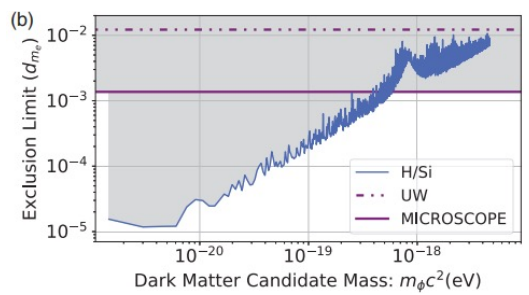
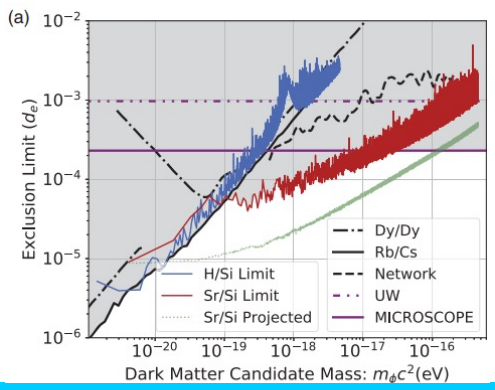
- Sr clock/cavity network comparison by GPS

Wcislo *et al*, Sci. Adv. 4, eaau4869 (2018)



- Sr clock/cavity/H maser comparisons

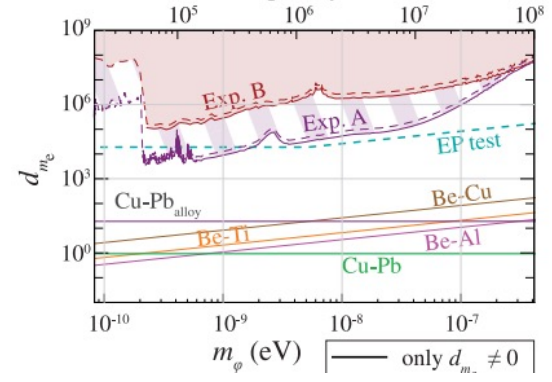
Kennedy *et al*, 125, 201302 (2020)



=> Spectroscopy experiments

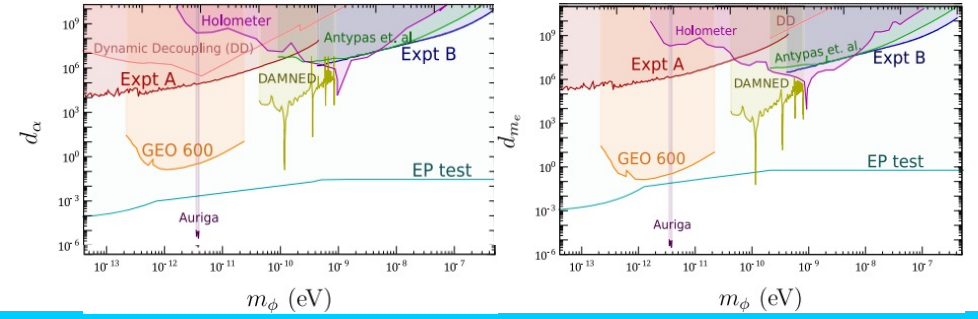
- Cs polarization spectroscopy

Tretiak *et al*, Phys. Rev. Lett. 129, 031301 (2022)



- I₂ Doppler/Doppler-free spectroscopy

Oswald *et al*, Phys. Rev. Lett. 129, 031302 (2022)



2) Potential from infrared precision measurements with $^{12}\text{C}_2\text{H}_2$

Acetylene molecular theory

Fundamental vibrational modes

	ν_1	3401.15 cm^{-1}	Symmetric CH stretching
	ν_2	1982.68 cm^{-1}	Symmetric CC stretching
	ν_3	3313.20 cm^{-1}	Antisymmetric CH stretching
	ν_4	608.99 cm^{-1}	Trans-bending
	ν_5	729.21 cm^{-1}	Cis-bending

Hermann, Mol. Phys. 105, 2217 (2007)

Polyads

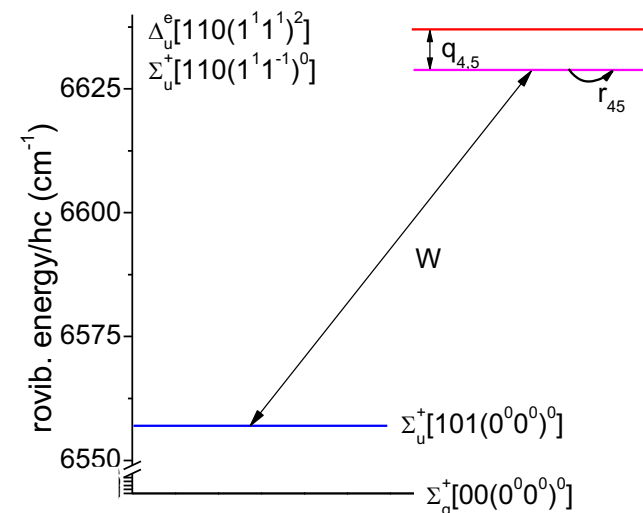
$$(\nu_1, \nu_2, \nu_3, \nu_4^{l_4}, \nu_5^{l_5})$$

$$k = l_4 + l_5; J$$

$$N_s = \nu_1 + \nu_2 + \nu_3$$

$$N_r = 5\nu_1 + 3\nu_2 + 5\nu_3 + \nu_4 + \nu_5$$

Energy levels ~196 THz



Effective Hamiltonian

Vibrational term: $G_v[v] = \sum_s \omega_s (\nu_s + g_s/2) + \sum_{s \leq s'} x_{ss'} (\nu_s + g_s/2)(\nu_{s'} + g_{s'}/2) + \sum_{k \leq k'} g_{kk'} l_k l_{k'}$

Rotational term: $F_r(J, k) = B[v](J(J+1) - k^2) - D[v](J(J+1) - k^2)^2 + H[v](J(J+1) - k^2)^3$

Hamiltonian matrix:

$$\begin{pmatrix} E[110(1^1 1^1)^0, J] + \frac{\rho_{45}}{4} J(J+1)(J(J+1)-2) & \frac{q_4 + q_5}{2} \sqrt{J(J+1)(J(J+1)-2)} & 0 \\ \frac{q_4 + q_5}{2} \sqrt{J(J+1)(J(J+1)-2)} & E[110(1^1 1^1)^0, J] + r_{45} & W \\ 0 & W & E[101(0^0 0^0)^0, J] \end{pmatrix}$$

Unperturbed energy levels:

$$E[110(1^1 1^1)^0, J] = (\nu_2 + g_{45}) + (B_2 + \gamma^{45})(J(J+1) - 4) - (D_2 + \delta^{45})(J(J+1) - 4)^2 + (H_2 + h^{45})(J(J+1) - 4)^3$$

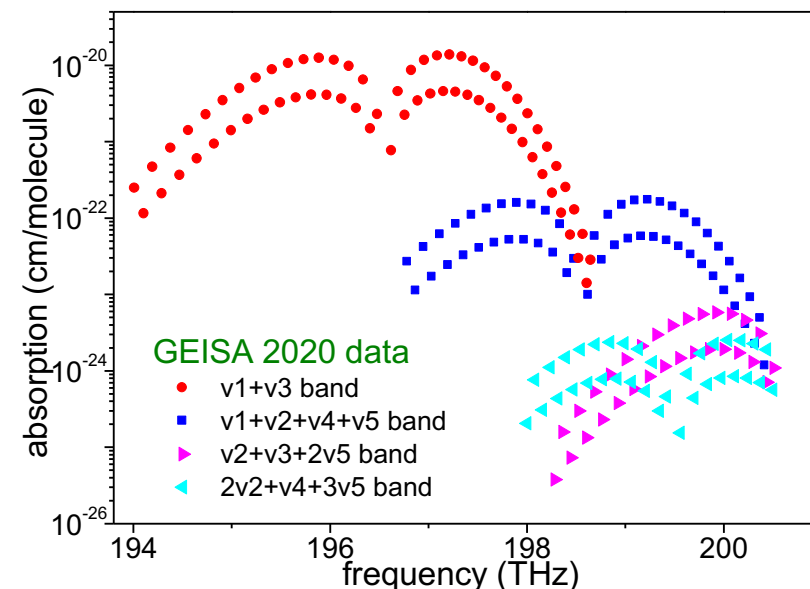
$$E[110(1^1 1^1)^0, J] = (\nu_2 - g_{45}) + (B_2 - \gamma^{45})(J(J+1) - 4) - (D_2 - \delta^{45})(J(J+1) - 4)^2 + (H_2 - h^{45})(J(J+1) - 4)^3$$

$$E[101(0^0 0^0)^0, J] = \nu_1 + B_1 J(J+1) - D_1 (J(J+1))^2 + H_1 (J(J+1))^3$$

$$E[000(0^0 0^0)^0, J] = B_0 J(J+1) - D_0 (J(J+1))^2 + H_0 (J(J+1))^3$$

Keppler *et al*, JMS 175, 411 (1996)

High-resolution spectroscopy



=> GEISA 2015 $^{12}\text{C}_2\text{H}_2$ linelist

Jacquinet-Husson *et al*, JMS 327, 31 (2016)

Sensitivity of the parameters to a variation of μ

Constant	Value (GHz)	K_A	Constant	Value (GHz)	K_A
B_0	35.274974565(42)	-0.9974937(4)	X_{14}	-416.1(42)	-1
$D_0 (\times 10^6)$	48.77824(39)	-2.01457(5)	X_{15}	-315.4(48)	-1
$H_0 (\times 10^{12})$	57(18)	-3	X_{22}	-223.4(21)	-1
$\alpha_1 (\times 10^3)$	206.986(33)	-1.5	X_{23}	-184.7(33)	-1
$\alpha_2 (\times 10^3)$	185.308(39)	-1.5	X_{24}	-381.6(30)	-1
$\alpha_3 (\times 10^3)$	176.332(33)	-1.5	X_{25}	-45.9(22)	-1
$\alpha_4 (\times 10^3)$	-40.5780(26)	-1.5	X_{33}	-828.3(36)	-1
$\alpha_5 (\times 10^3)$	-66.9159(12)	-1.5	X_{34}	-300.4(42)	-1
$\gamma^{44} (\times 10^4)$	-19.720(33)	-2	X_{35}	-280.0(48)	-1
$\gamma^{55} (\times 10^4)$	-32.954(15)	-2	X_{44}	103.7(16)	-1
$\gamma^{45} (\times 10^4)$	-67.641(12)	-2	X_{45}	-66.9(30)	-1
$\beta_1 (\times 10^8)$	-42.45(27)	-2.5	X_{55}	-71.2(16)	-1
$\beta_2 (\times 10^8)$	5.93(44)	-2.5	g_{44}	23.42939(27)	-1
$\beta_3 (\times 10^8)$	-40.99(40)	-2.5	g_{45}	198.00315(24)	-1
$\beta_4 (\times 10^8)$	103.13(23)	-2.5	g_{55}	104.21008(19)	-1
$\beta_5 (\times 10^8)$	77.76(18)	-2.5	r_{45}	-187.03212(33)	-1
$\delta^{44} (\times 10^{10})$	-517(18)	-3	$r_{J45} (\times 10^4)$	58.565(19)	-2
$\delta^{45} (\times 10^{10})$	-1184(22)	-3	$\rho_{45} (\times 10^8)$	-52.79(90)	-2
$\delta^{55} (\times 10^{10})$	-446(17)	-3	$q_{04} (\times 10^3)$	-157.3485(36)	-1.5
B_1	34.8877668(57)	-0.991972(1)	$q_{05} (\times 10^3)$	-139.7165(36)	-1.5
$D_1 (\times 10^6)$	47.8852(39)	-2.0061(2)	$q_{44} (\times 10^5)$	53.45(33)	-2
B_2	35.002840(20)	-0.993256(1)	$q_{45} (\times 10^5)$	-237.3(1.3)	-2
$D_2 (\times 10^6)$	48.107(60)	-2.0277(36)	$q_{54} (\times 10^5)$	-330.1(1.4)	-2
v_1	196576.8494(18)	-0.47159(9)	$q_{55} (\times 10^5)$	-113.86(36)	-2
v_2	198922.7020(21)	-0.48114(7)	$q_{J4} (\times 10^8)$	117.82(19)	-2.5
X_{11}	-743.8(66)	-1	$q_{J5} (\times 10^8)$	115.351(69)	-2.5
X_{12}	-350.2(39)	-1	W	193.87198(96)	-1
X_{13}	-3223(16)	-1	$W_J (\times 10^5)$	-101.93(60)	-2

sensitivity coefficient

$$K_A = d \ln A / d \ln \mu$$

Dependence on molecular structure parameters => sensitivity to μ -variation

Constantin, Vibrational Spectr. 85, 228 (2016)

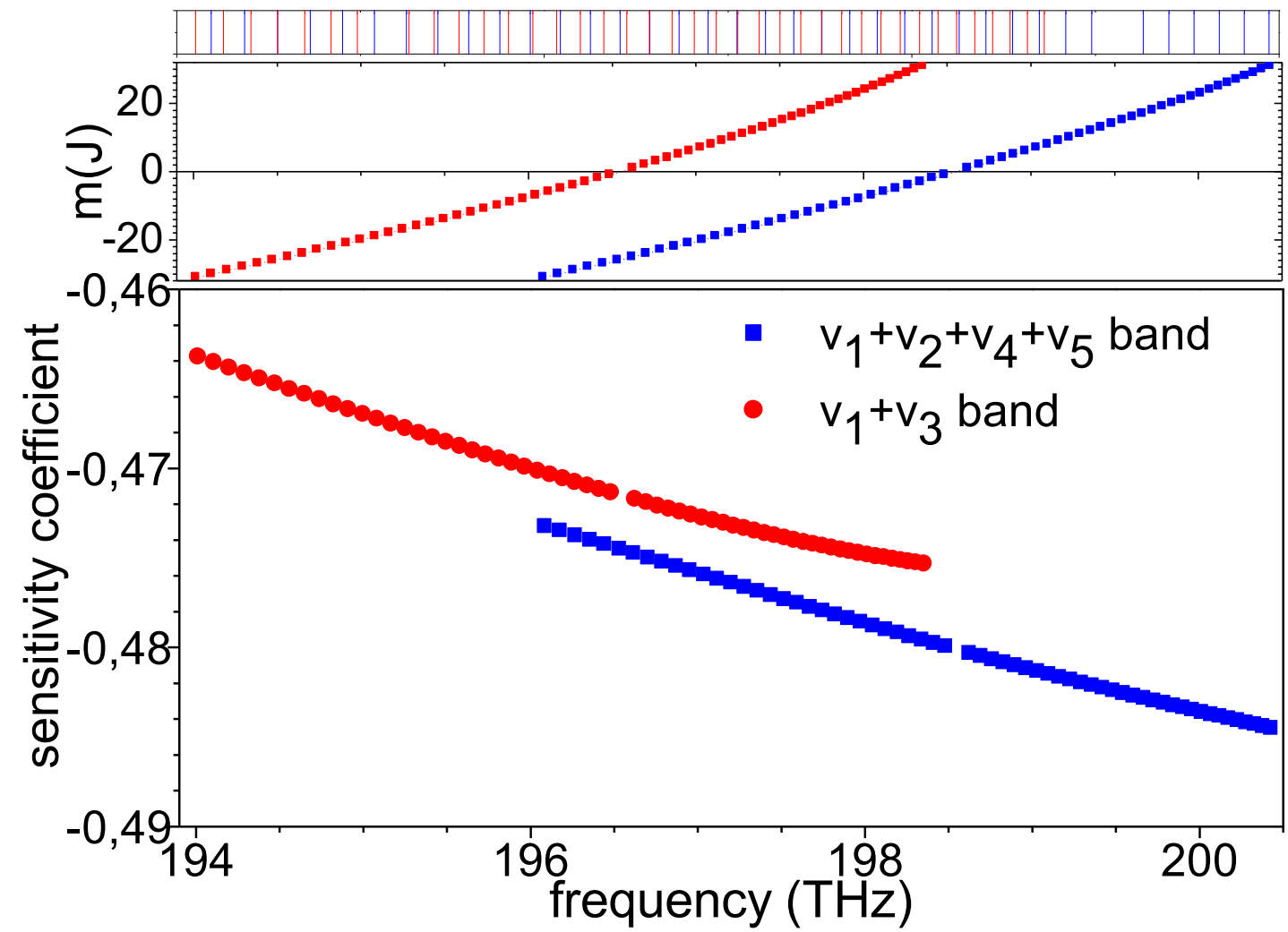
• Sensitivity of the transitions to a variation of μ

-Hamiltonian approach to calculate sensitivity to μ -variation of $^{12}\text{C}_2\text{H}_2$ reference transitions

Constantin, Vibrational Spectr. 85, 228 (2016)

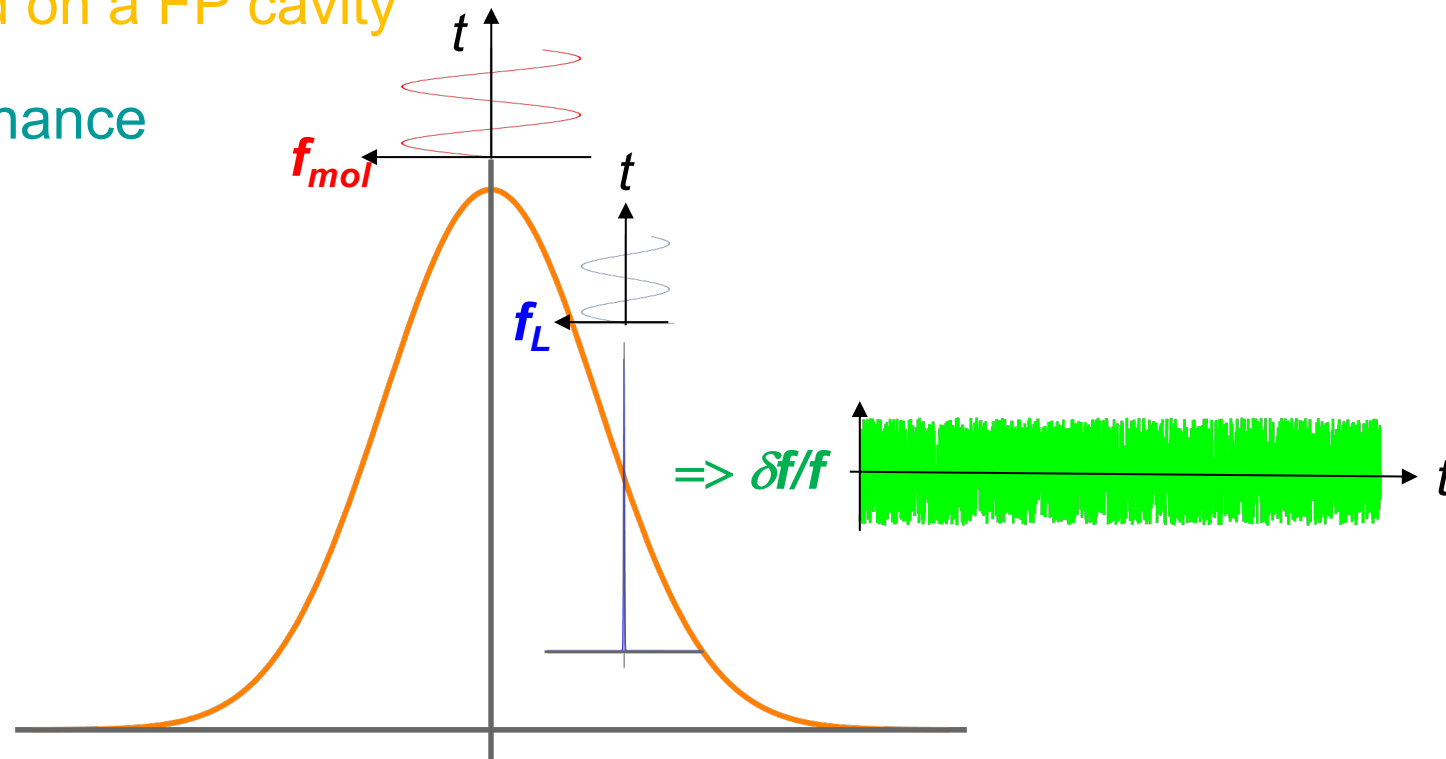
Prediction of acetylene frequency in function of μ & derivative of simulated data

=>Sensitivity of reference transition P(16) of the $\nu_1+\nu_3$ band of $^{12}\text{C}_2\text{H}_2$: $K_\mu=-0.468$



- Principle of the experiment

- Laser stabilized on a FP cavity
- Molecular resonance



$$\begin{aligned}
 \frac{f_L(t) - f_{mol}(t)}{f_L(t)} = & \left(Q_\alpha^L h^L(f_\phi) - Q_\alpha^{mol} h^{mol}(f_\phi) \right) \frac{\Delta\alpha(t)}{\alpha} \\
 & + \left(Q_\mu^L h^L(f_\phi) - Q_\mu^{mol} h^{mol}(f_\phi) \right) \frac{\Delta\mu(t)}{\mu}
 \end{aligned}$$

- Sensitivity coefficients : $Q_\alpha^L \cong 1$; $Q_\mu^L \cong 0$; $Q_\alpha^{mol} = 2$; $Q_\mu^{mol} = -0.47$

Pařteka *et al*, PRL 122, 160801 (2019); Constantin, Vibrational Spectr. 85, 228 (2016)

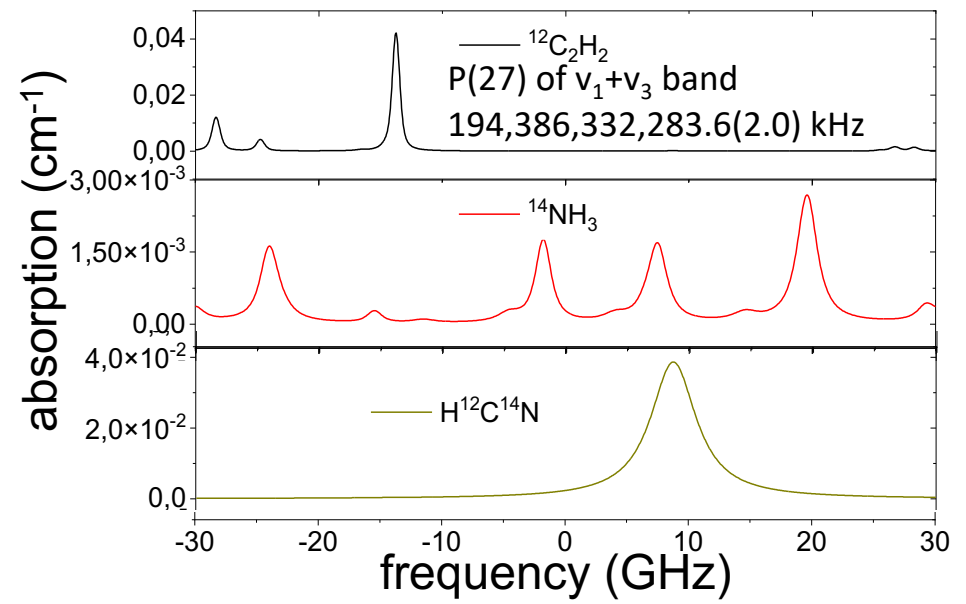
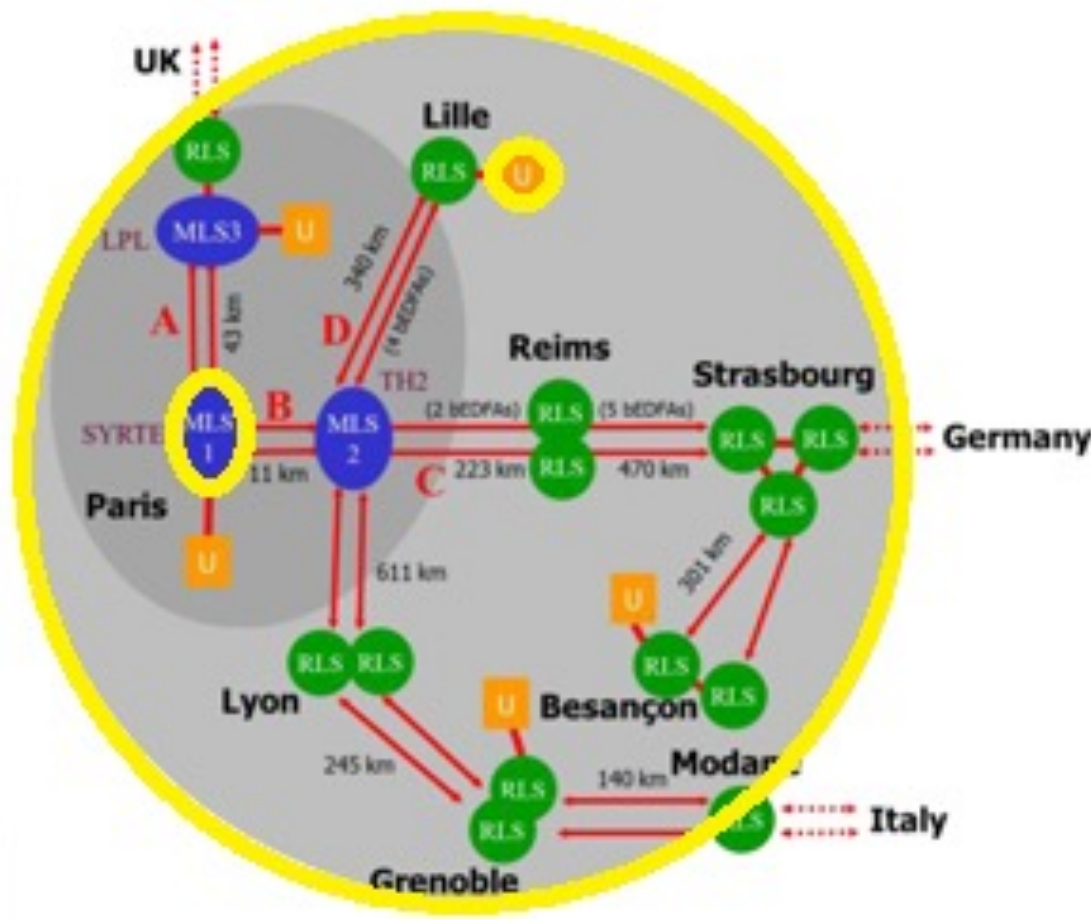
- Response functions h^L , h^{mol} with high-frequency cutoffs

-frequency cutoffs : linewidth of the molecular line, delay in propagation of sound in the ULE spacer, ...

• Towards a network spectrometer

- REFIMEVE network :
 - 1542 nm laser locked to ULE cavity/H maser
 - phase-stabilized fiber link
- Cantin *et al*, New J. Phys. 23, 053027 (2021)

=> Signal to PhLAM almost continuously at 194,400,084,500.000(25) kHz



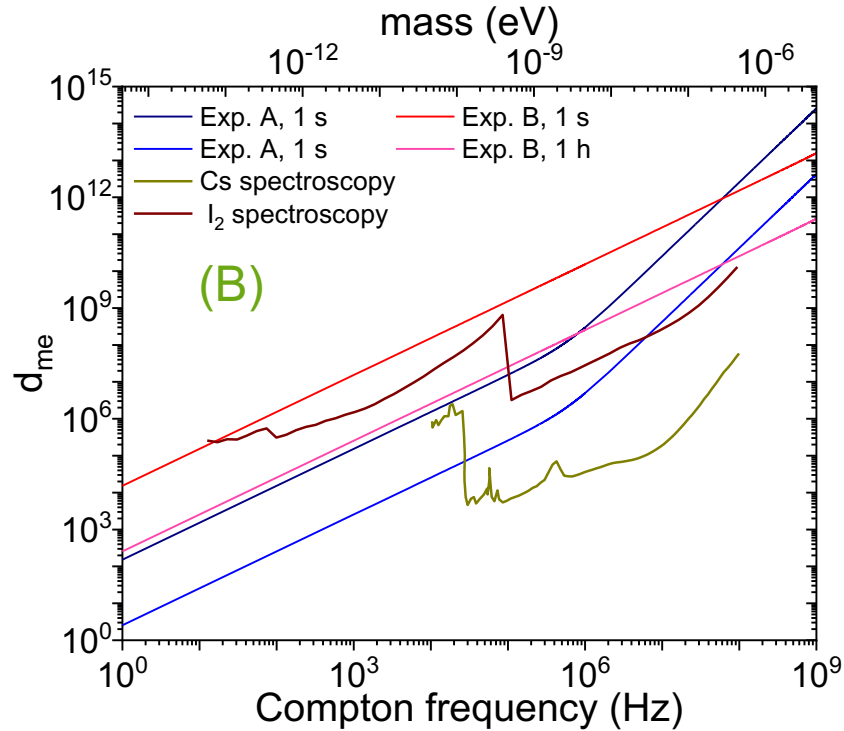
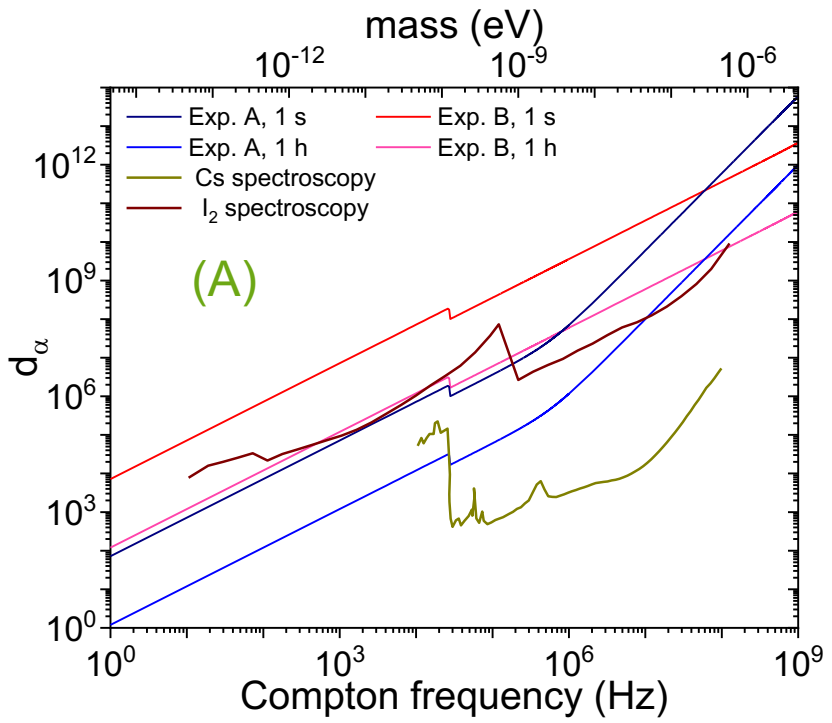
- Broadband tuning with optical modulators -LiNbO₃ photonic platform
- Fast data acquisition at 1 Gsa/s level -fast DAQ and data storage
- Low-noise absorption detection -differential detection, noise-eater implementation,...

=> fractional stability $< 10^{-15} (\tau/1\text{s})^{-1/2}$
 => fractional uncertainty $< 10^{-14}$

=> compact/robust molecular spectrometer
 => integration into the optical fiber network
 Constantin, Proc. IFCS-EFTF 2023 Paper Id 7354

• Estimated sensitivity of the experiment

- Exp. A : recording a Doppler-free molecular line with $\delta f/f=10^{-14} (\tau/1s)^{-1/2}$
- Exp. B : recording a linear absorption molecular line with $\delta f/f=10^{-12} (\tau/1s)^{-1/2}$
- High-frequency cutoffs in the response of the experimental setup
 - Sound propagation in ULE spacer $f_{c1}=27$ kHz
 - Molecular transition sub-Doppler $f_{c2A}=100$ kHz and linear absorption $f_{c2B}=4$ GHz linewidths



=>improved limits on Compton frequencies domain by one order of magnitude

=>potential to improve the state-of-the-art constraints at low frequencies by averaging

2) Potential from frequency measurements of a MW acetylene clock

• Principle of the proposed experiment

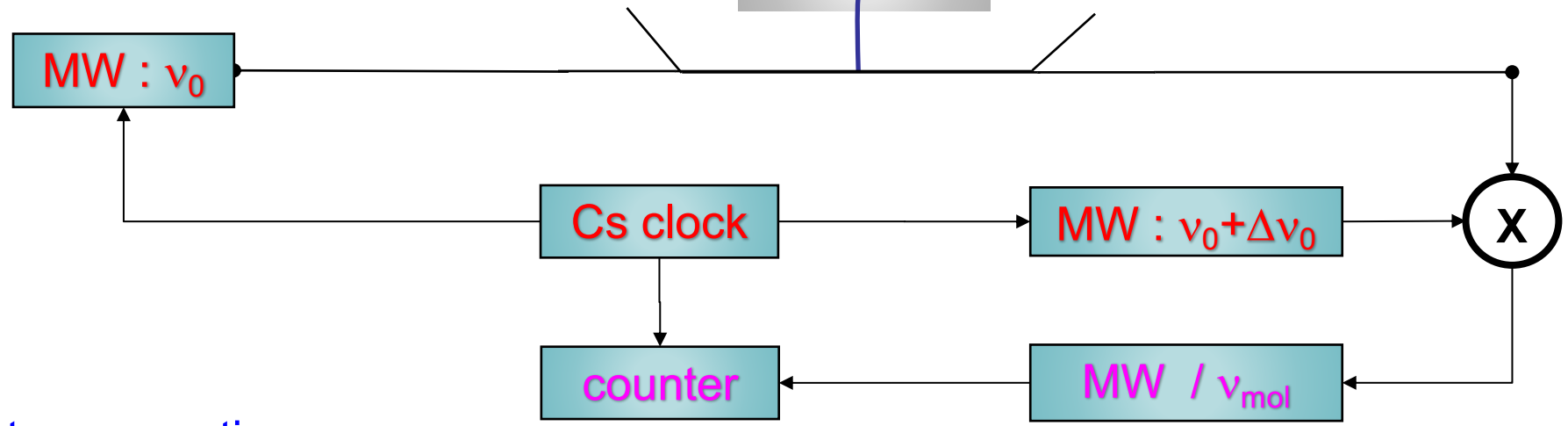
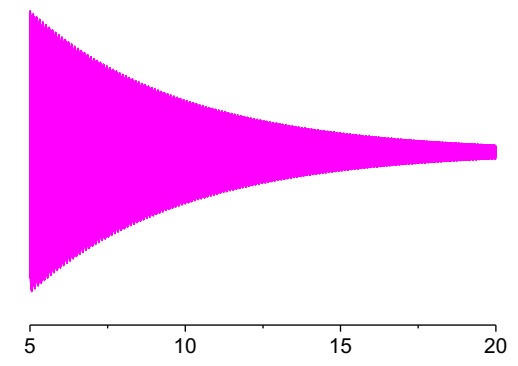
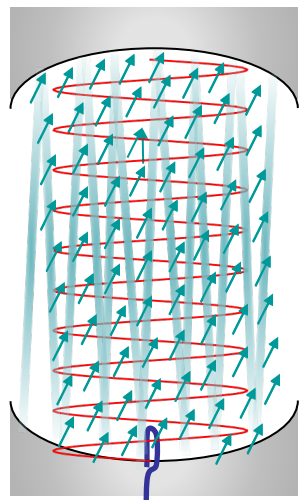
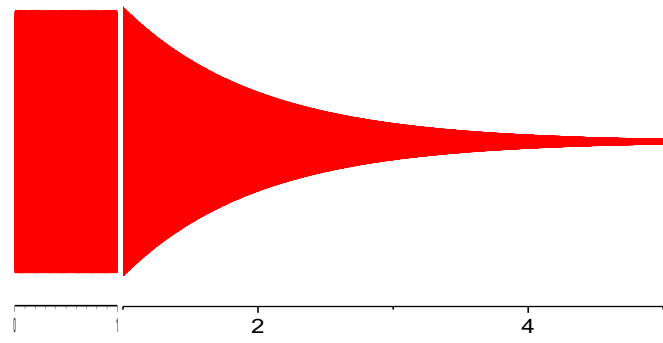
Measurement of a MW transition with enhanced sensitivity coefficient to μ -variation

Constantin, Proc. CLEO-Europe/EQEC paper ED-1.4 (2023)

$$E_{cav} = E_0 \cdot \cos(2\pi\nu_0 t) \text{ for } 0 < t < \tau_p;$$

$$E_0 \cdot e^{-(t-\tau_p)/\tau_{cav}} \cos(2\pi\nu_0(t - \tau_p)) \text{ for } t \geq \tau_p$$

$$E_{FID} = \exp\left(-2\pi\Delta\nu_{HWHM}^{press} t - (\pi\Delta\nu_{HWHM}^{Doppler} t)^2 / \ln(2)\right) \times E_{mol} \cdot \cos(2\pi\nu_{mol}(t - \tau_p - 4\tau_{cav}) + \theta)$$



State preparation
 -heating
 -optical pumping
 -MW $\pi/2$ pulse

Cavity ringdown
 Molecular FID

Heterodyne mixing
 FFT+MW osc. locking

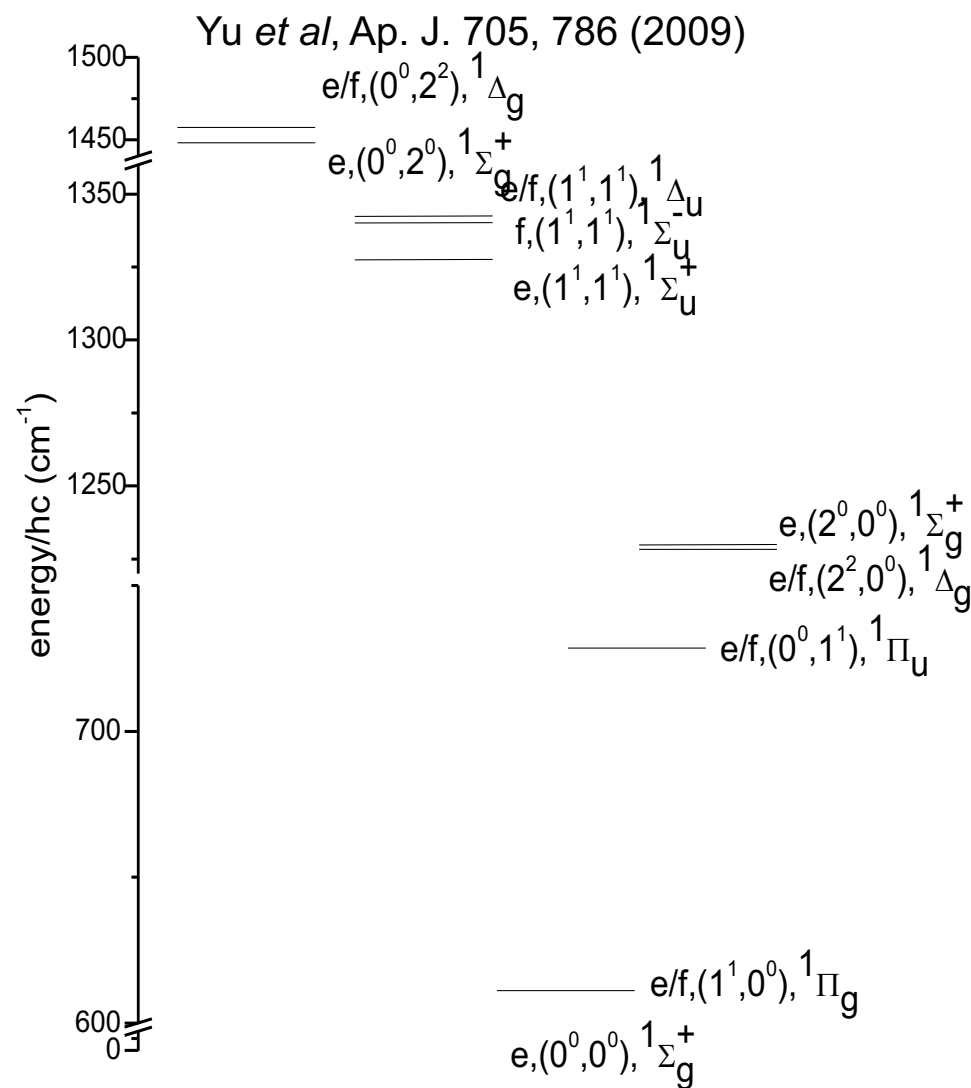
C_2H_2/Cs frequency measurement

• $^{12}\text{C}_2\text{H}_2$ low energy levels and rovibrational interactions

Vibrational energy: $v_4^{l_4}, v_5^{l_5}$

$$E_V/hc = \sum_{i=4}^5 \omega_i \left(v_i + \frac{d_i}{2} \right) + \sum_{i,j=4}^5 x_{ij} \left(v_i + \frac{d_i}{2} \right) \left(v_j + \frac{d_j}{2} \right) + \sum_{i,j=4}^5 g_{ij} \left(l_i + \frac{d_i}{2} \right) \left(l_j + \frac{d_j}{2} \right) + \dots$$

$\omega_i \sim 1/\sqrt{\mu}; x_{ij} \sim 1/\mu; g_{ij} \sim 1/\mu, \dots$



Vibrational energy: $v_4^{l_4}, v_5^{l_5}$

$$E_V/hc = \sum_{i=4}^5 \omega_i \left(v_i + \frac{d_i}{2} \right) + \sum_{i,j=4}^5 x_{ij} \left(v_i + \frac{d_i}{2} \right) \left(v_j + \frac{d_j}{2} \right) + \sum_{i,j=4}^5 g_{ij} \left(l_i + \frac{d_i}{2} \right) \left(l_j + \frac{d_j}{2} \right) + \dots$$

$\omega_i \sim 1/\sqrt{\mu}; x_{ij} \sim 1/\mu; g_{ij} \sim 1/\mu, \dots$

Rotational energy: $J, k=l_4+l_5$

$$E_R/hc = B(v_4^{l_4}, v_5^{l_5}) [J(J+1) - k^2] + D(v_4^{l_4}, v_5^{l_5}) [J(J+1) - k^2]^2 + \dots$$

$$B(v_4^{l_4}, v_5^{l_5}) = B_0 - \sum_{i=4}^5 \alpha_i \left(v_i + \frac{d_i}{2} \right) + \dots$$

$$D(v_4^{l_4}, v_5^{l_5}) = D_0 + \sum_{i=4}^5 \beta_i \left(v_i + \frac{d_i}{2} \right) + \dots$$

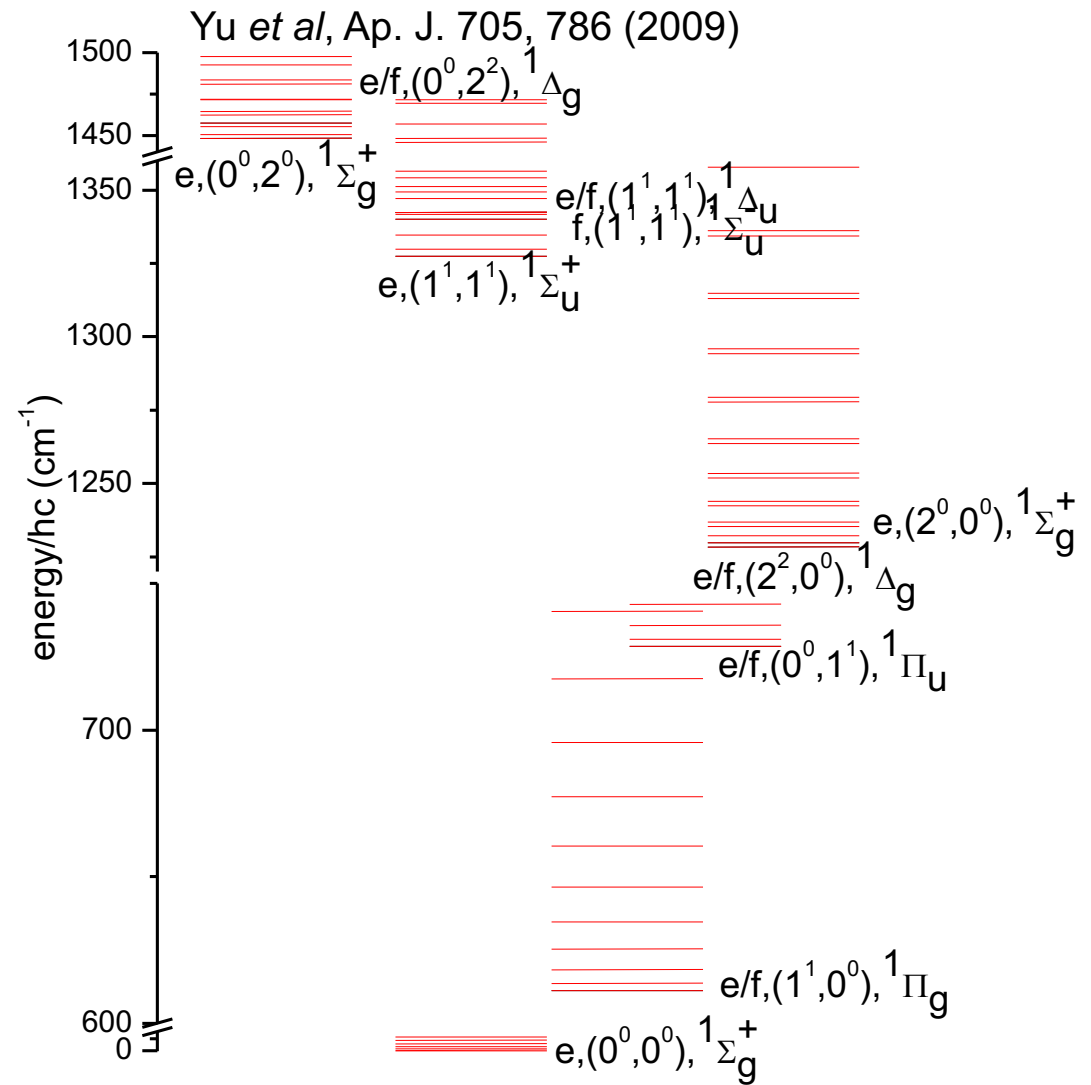
$B \sim 1/\mu; D \sim 1/\mu^2; \alpha \sim 1/\mu^{3/2}; \beta \sim 1/\mu^{5/2}$

Rotational & vibrational I-type interactions and doublings

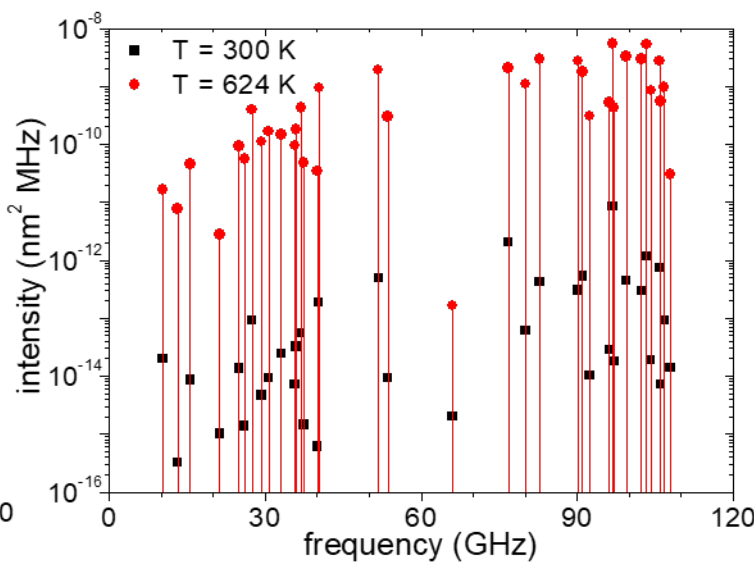
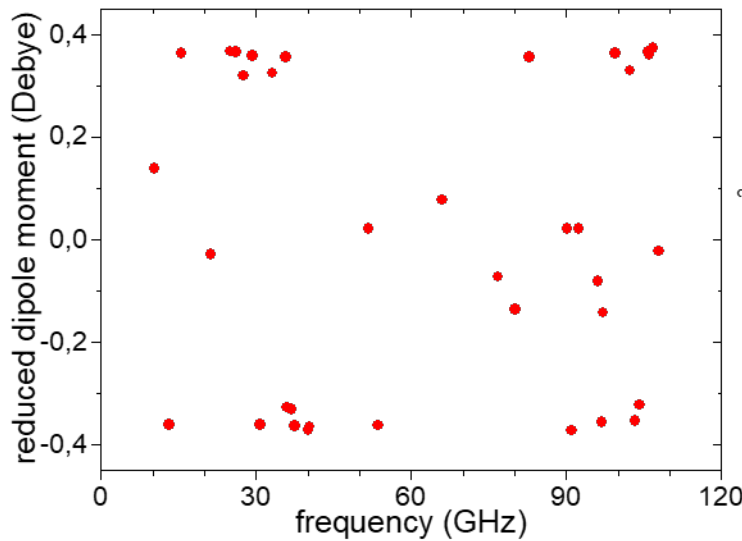
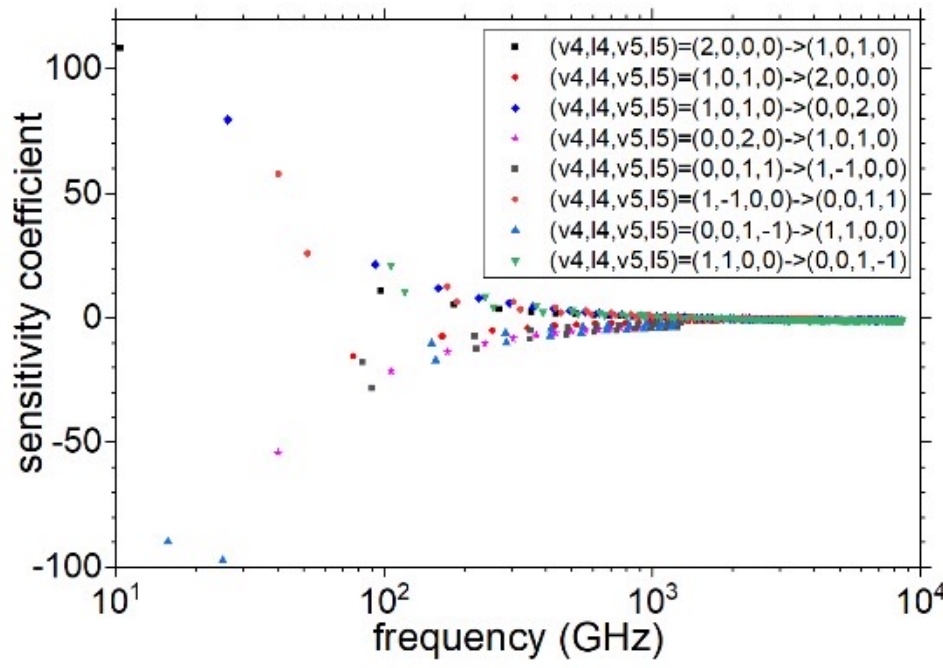
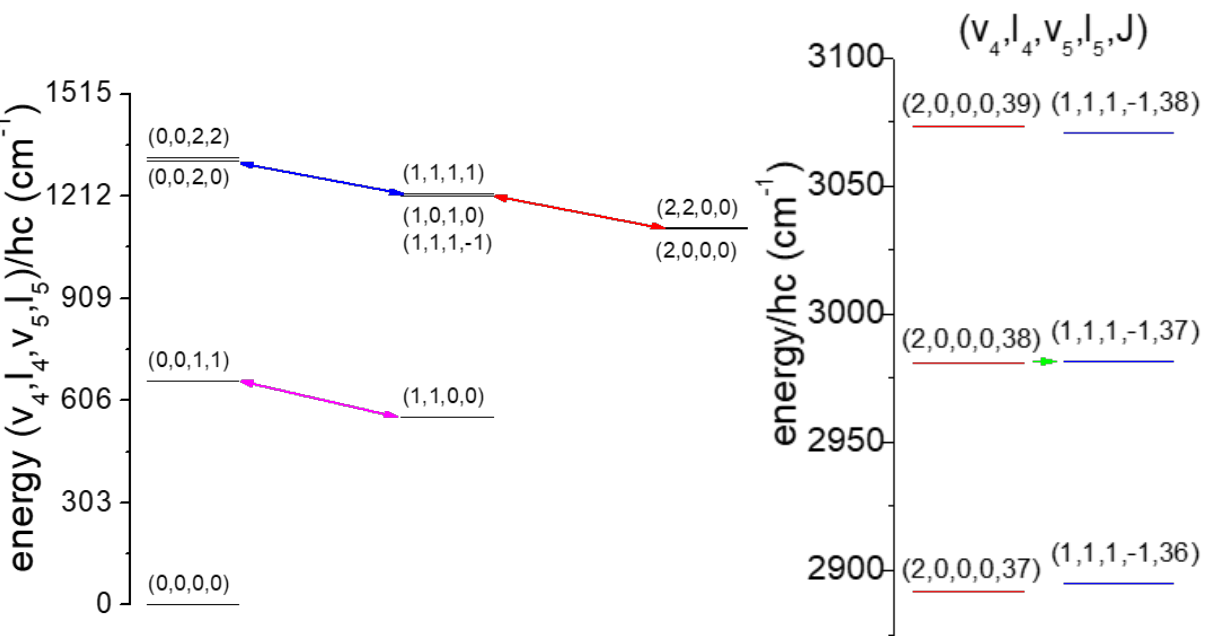
For exemple $M = J(J+1) - k^2$:

$$\begin{pmatrix} E_R(J, 0) + E_V(2^{0,\pm 2}, 0^0) & (\sqrt{2} hcq_4(2,0)/2) \sqrt{M(M-2)} \\ 0 & E_R(J, 0) + E_V(2^{0,\pm 2}, 0^0) \pm (hpc_4/2) M(M-2) \end{pmatrix}$$

=> Energy level prediction with 10^3 - 10^7 Hz accuracy ; estimation of the μ -sensitivity



•MW transitions with enhanced sensitivity to μ -variation



• Frequency stability from a $^{12}\text{C}_2\text{H}_2$ microwave line

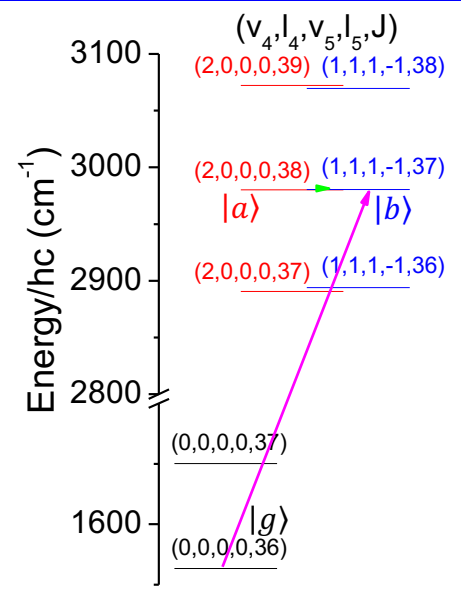
Transition $(v_4^{l_4}, v_5^{l_5}, J, \text{sym}) = (2^0, 0^0, 38, e^1 \Sigma_g^+) \rightarrow (1^1, 1^{-1}, 37, e^1 \Sigma_u^+)$

Exp. param. : $v_{\text{mol}} = 10363 \text{ MHz}$, $L = 1 \text{ m}$, $w_0 = 0.1 \text{ m}$, $Q_0 = 10^5$
 $\pi/2 \text{ MW pulse}$; full population transfer from $|g\rangle$ to $|b\rangle$

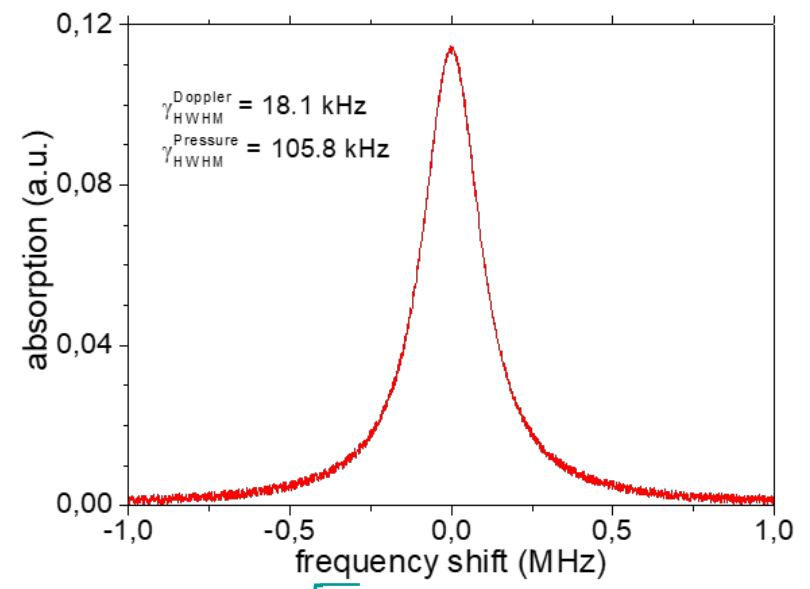
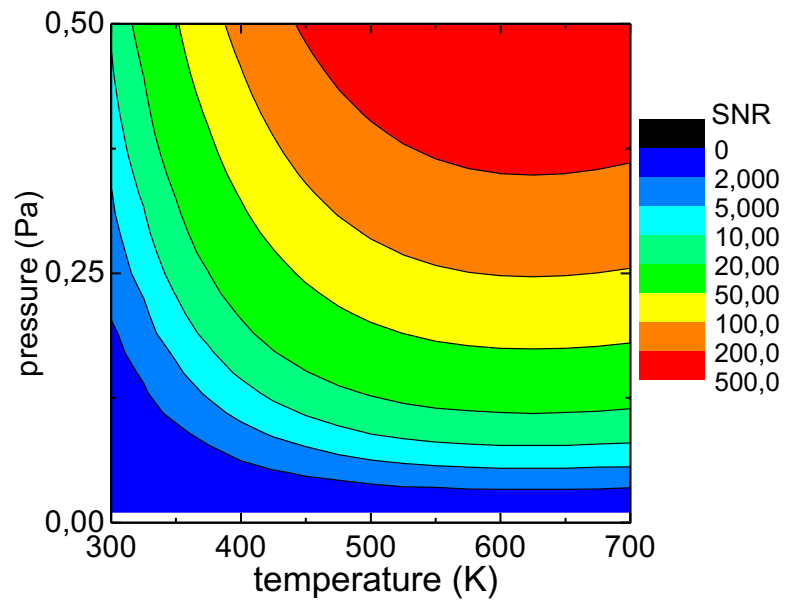
FID power :
$$P_{\text{FID}} = \frac{16}{9} \pi^2 Q_0 v_0 (\mu_{ab} \Delta N_0)^2 \pi w_0^2 L \left(\frac{\int_0^{\kappa E_0 \tau_p} J_1(u) du}{\kappa E_0 \tau_p} \right)^2$$

Campbell *et al*, J. Chem. Phys. 74, 813 (1980)

Noise power :
$$P_{\text{noise}} = k T_{\text{noise}} B \approx k T_{\text{exp}} / \tau_{\text{cav}}$$



Optimization with constraint $B = \gamma_{\text{cav}} \sim \gamma_{\text{mol}} \Rightarrow T_{\text{exp}} = 624 \text{ K}$; $P_{\text{exp}} = 0.24 \text{ Pa}$; $\text{SNR} = 94$



$$\Rightarrow \text{Allan fractional instability} : \sigma_{\delta f_{\text{mol}}/f_{\text{mol}}}(\tau) \approx \frac{P_{\text{FID}}/P_{\text{noise}}}{Q_{\text{mol}}} \sqrt{\frac{T_c}{\tau}} \approx \frac{2.1 \times 10^{-10}}{\sqrt{\tau/1s}}$$

• Systematic frequency shifts of a $^{12}\text{C}_2\text{H}_2$ microwave line

Transition $(v_4^{l_4}, v_5^{l_5}, J, \text{sym}) = (2^0, 0^0, 38, e^1 \Sigma_g^+) \rightarrow (1^1, 1^{-1}, 37, e^1 \Sigma_u^+)$ at 10363 MHz

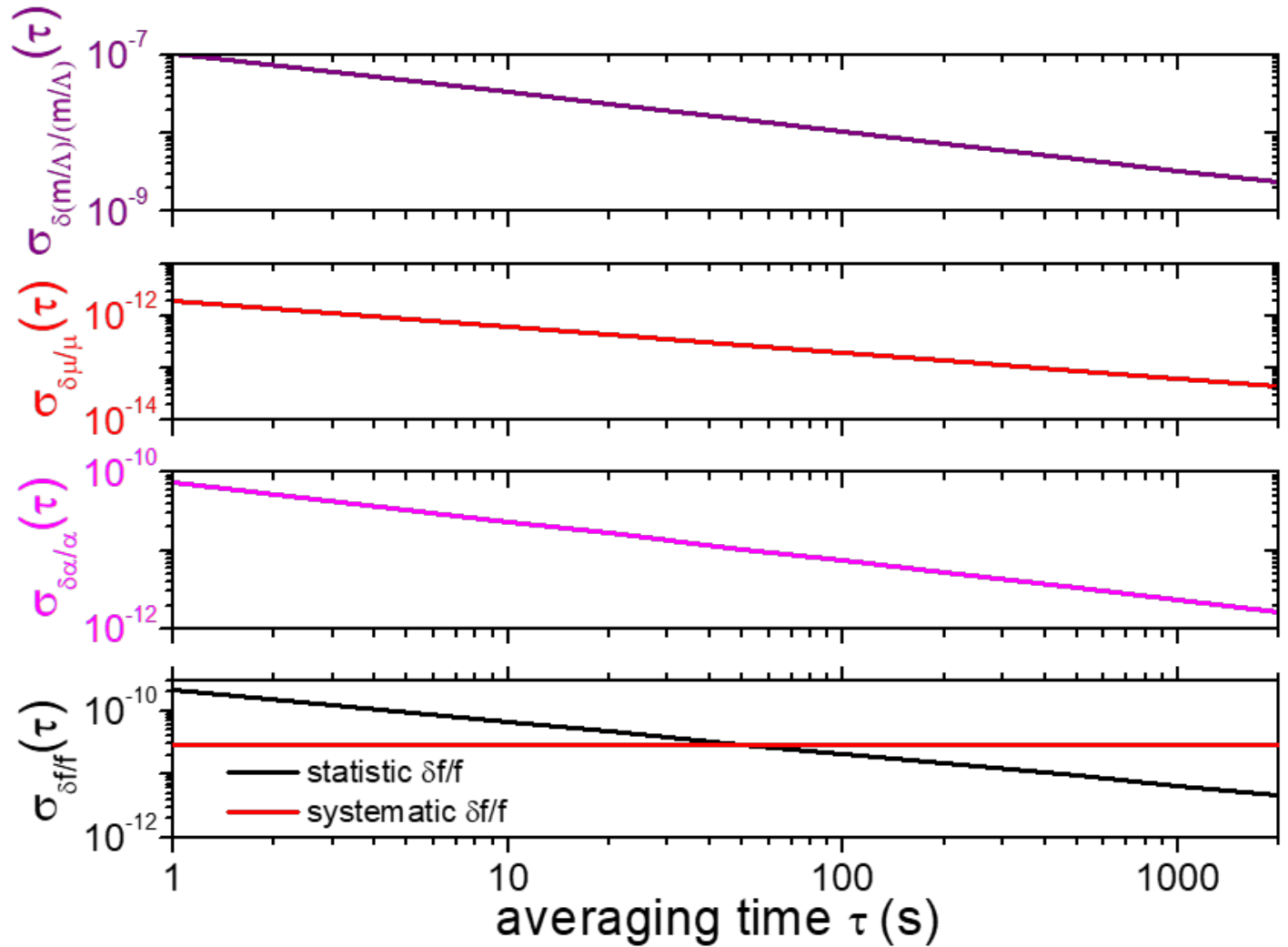
Systematic effect	value	unc.
Cavity pulling	44 kHz	0.22 Hz
$f_{mol} - f_{mol,0} = (f_{cav} - f_{mol}) \times (Q_{cav}/Q_{mol})^2$		
param.:	$\alpha_T = 5 \times 10^{-6}/^\circ\text{C}$	
uncertainty:	$\Delta T = 1^\circ\text{C}$ $(f_{cav} - f_{mol}) = \Delta\nu/10$	
Viennet <i>et al</i> , IEEE TIM 21, 204 (1972)		
Pressure shift	246 Hz	0.19 Hz
extrapolation meas. v_4+v_5 band at high T and low p		
param.:	T=624 K; p=0.24 Pa; no J-dependence	
uncertainty:	$\Delta T = 1\text{ K}$	
Dhyne <i>et al</i> , JQSRT 112, 969 (2011)		
Zeeman shift	28 Hz	86 mHz
extrapolation meas. MVCD v_4 & v_5 band; eval. M=J states		
param.:	B=48 μT	
uncertainty:	$\Delta B = 0.15 \mu\text{T}$	
Tam <i>et al</i> , J. Chem. Phys. 104, 1813 (1996)		

Systematic effect	value	unc.
DC Stark shift	6.4 μHz	1.3 μHz
1st order perturbation theory, eval. for M=0 states		
param.:	E=1 V/m	
uncertainty:	$\Delta\mu/\mu = 20\%$	
Barnes <i>et al</i> , Chem. Phys. Lett. 237, 437 (1995)		
BBR shift	1.4 mHz	9 μHz
ab-initio scalar polarizabilities, eval. for M=J states		
param.:	T=624 K	
uncertainty:	$\Delta T = 1\text{ K}$	
Russell and Spackman, Mol. Phys. 88, 1109 (1996)		
SODE shift	-34.3 mHz	55 μHz
param.:	T=624 K	
uncertainty:	$\Delta T = 1\text{ K}$	
Recoil doubling	negligible	
Autler-Townes doubling	negligible	

• Sensitivity to variations of fundamental constants

$$\frac{\Delta(f_{MW/mol}/f_{Cs})}{f_{MW/mol}/f_{Cs}} = (Q_{\alpha}^{mol} - Q_{\alpha}^{Cs}) \frac{\Delta\alpha}{\alpha} + (Q_{\mu}^{mol} - Q_{\mu}^{Cs}) \frac{\Delta\mu}{\mu} + (Q_q^{mol} - Q_q^{Cs}) \frac{\Delta(\hat{m}/\Lambda)}{\hat{m}/\Lambda}$$

- Sensitivity coefficients : $Q_{\alpha}^{Cs} = 2.83$; $Q_{\mu}^{Cs} = -1$; $Q_q^{Cs} = 0.002$; $Q_{\alpha}^{mol} = 0$; $Q_{\mu}^{mol} = 109$; $Q_q^{mol} = 0$



- Accurate theory for modeling frequency & frequency shifts of $^{12}\text{C}_2\text{H}_2$ transitions
- Accurate determination of sensitivity coefficients to μ -variation
- Enhanced sensitivity coefficients for MW transitions between near-resonant levels
- Development of spectrometers at 194 THz using the REFIMEVE network with precision at levels of $(\delta f/f)^{\text{Doppler-free}} = 10^{-14} (\tau/1\text{s})^{-1/2}$ and $(\delta f/f)^{\text{Doppler}} = 10^{-12} (\tau/1\text{s})^{-1/2}$
- Microwave molecular clock with $(\delta f/f)^{\text{A}} = 2.1 \times 10^{-10} (\tau/1\text{s})^{-1/2}$ and $(\delta f/f)^{\text{B}} = 2.9 \times 10^{-11}$
- Potential to improve constraints on UDM couplings to SM particles

Thank you for your attention !

Acknowledgements

