







# Searching for UltraLight Dark Matter with laboratory experiments

#### J. Gué, <u>A. Hees</u>, P. Delva, P. Wolf

LNE-SYRTE, CNRS, Paris Obs., Université PSL, Sorbonne Univ.



GC observation @Keck





Systèmes de Référence Temps-Espace

#### in collaboration with

- SYRTE: M. Abgrall, S. Bize, E. Cantin, F. Florian, R. Le Targat, J. Lodewyck, P-E. Pottie, ...
- CEA: P. Brun, L. Chevalier, H. Deschamps, P. Polovodov, E. Savalle
- CSM and OCA: O. Minazzoli
- OCA/Grasse Station: J. Chabé, C. Courde
- ROB: B. Bertrand, P. Defraigne
- U. Sidney: Y. Stadnik
- U. Queensland: B. Roberts

#### DM needed to explain astro/cosmo observations but not direct detection so far

• DM needed at: galactic scales (rotation curves, ...), galaxy cluster (bullet cluster, ...), cosmo (CMB, structure formation, ...)



# UltraLight Dark Matter needs to be a boson and it behaves classically

• Occupation number (number of particles per volume of phase-space)

$$\frac{n}{n_k} \sim \frac{6\pi^2 \hbar^3 \rho_{\rm DM}}{m^4 c^2 v_{\rm max}^3}$$

Calculation inspired from Tourrenc et al, arXiv:quantum-ph/0407187, 2004

- In our Galaxy  $\rho_{\rm DM} \approx 0.4 {\rm GeV/cm}^3$
- This occupation number is larger than 1 if the DM mass is lower than
  ~ 10 eV: Dark Matter lighter than 10 eV can only be made of boson
  - a bosonic scalar particle (i.e. a scalar field)
  - a bosonic pseudo-scalar particle (i.e. an axion)
  - a boson vector particle (i.e. a hidden photon)
- For m << eV: the occupation number is huge and such a bosonic field can be treated classically (no quantization)

## Mainly two phenomenological signatures explored so far

I. Oscillatory behaviour of the additional field

see Arvanitaki et al, PRD, 2015

- oscillation with stochastic amplitude
- oscillation with amplitude depending on location (screening/ scalarization possible)

2. Topological defect: domain wall, ...

see Roberts et al, Nature Com., 2017

- search for transients signatures in the data
- Search using fiber-link comparison of clocks or using GNSS Galileo data (+ dedicated SLR campaign)

see Roberts et al, New. Journal of Phys. 2020 Bertrand et al, submitted to ASR, 2023

## A massive scalar field or a massive vector field oscillates at its Compton frequency

• A massive scalar field  $\varphi$  |• A massive vector field  $X_{\mu}$ 

• When  $H \ll m^2$  (H=Hubble constant):



## If the new field makes DM, its oscillation amplitude is related to the DM energy density

- A massive scalar field  $\varphi$  |• A massive vector field  $X_{\mu}$
- Oscillates at Compton frequency

$$\varphi = \varphi_0 \cos mt \qquad \qquad \vec{X} = \vec{X}_0 \cos mt$$

• The averaged stress-energy tensor

 $\rho \sim \left\langle T_0^0 \right\rangle = \frac{m^2 \varphi_0^2}{2}$ 

$$\rho = \frac{m^2 \left| \vec{X}_0 \right|^2}{2}$$

$$p_{ij} \sim \left\langle T_j^i \right\rangle = 0$$

• The scalar/vector field can be identified as a pressureless fluid

 $\Rightarrow$  a possible Dark matter candidate!

## In experimental searches, we look for interactions between this new fields and SM

- Different couplings for different fields:
  - scalar: dilaton couplings  $d_i$  (to EM, fermions, QCD): constants of Nature ( $\alpha$ , fermion masses) depend on space/time [atomic clocks, UFF experiments, ...]
  - Axion (pseudo-scalar): coupling to pseudo scalar Lagrangian density (EM, QCD, fermion). Recent result: mass of pions depend quadratically to the axion field [UFF violation]

see Kim and Perez, arXiv 2023

 Vector/Dark photon: kinetic mixing to EM χ [modification of EM, ...], coupling to the fermonic currents [B, B-L couplings, leading to a UFF violation]
 See Horns, JCAP, 2011 P. Fayet, PRD, 2018

### Vector Ultra Light Dark Matter: A Dark Photon/Hidden Photon

# A vector DM will interact with electromagnetism

• An effective Lagrangian for the vector-matter coupling

$$\mathcal{L}_{\text{mat}}\left[\Psi, g_{\mu\nu}, X_{\mu}\right] = \mathcal{L}_{\text{SM}}\left[\Psi, g_{\mu\nu}\right] - \frac{\chi}{2} F^{\mu\nu} X_{\mu\nu} + \dots$$

see Horns et al, JCAP, 2013 and references therein

- Kinetic mixing coupling  $\chi$  characterises the coupling with EM
- Other couplings with matter can be considered like to the B-L current: leads to a violation of the UFF

A hidden photon field will generate a small EM field and vice versa

### An oscillating DM vector field will generate a small electric field

- Oscillating DM vector field  $\vec{X} = \vec{X}_0 \cos mt$  will generate an EM field  $\vec{E}_{\rm DM} = -\partial_t \vec{A} = -m\chi \vec{X}_0 \sin mt$  $\rho = \frac{m^2 \left| \vec{X_0} \right|^2}{2}$
- Reminder: the amplitude is related to the DM density
- Idea to search for such a DM: amplify this electric field using reflectors (boundary condition: creation of a classical propagating EM field)



#### Use cavity and Rydberg spectro to search for Dark Photon



- 2 electric fields: (i) an injected field and (ii) the DM induced field
- Detection of  $E^2$  through the Stark effect  $\Delta\nu\propto E^2$  using Rydberg atoms

$$E^2 \in \chi \overrightarrow{E_{in}} \cdot \overrightarrow{X_{DM}} \cos\left(\omega_{DM} - \omega_{in}\right) t + \dots$$

- Slowly evolving signal (detectable)

Enhance the signal amplitude

- Large range of DM mass explorable

#### Sensitivity analysis

#### work from J. Gué, PhD student

 Analysis includes: cavity losses (Q-factor), statistical noise (Rydberg), systematic noise (RIN of injected field), ...



New proposal of experiment: cavity can be used to search for Dark Photons (acts as a narrow band resonant detector)



• For a spherical dish, the electric field will be focused at the center + non-relevant electric field will be focused at the focal point

#### First result: SHUKET puts a stringent constraint on the kinetic mixing parameter



from P. Brun et al, PRL, 2019

#### Recent update from CEA-SYRTE collaboration

#### work from J. Gué, PhD student

- Improved data analysis considering signal stochasticity: improves slightly the constraints
   based on a methodology presented in E. Savalle et al, PRL, 2021
- Improved modelling of the experiment including
  - Diffraction of the EM field emitted by the dish
  - matching with the EM mode of the antenna

This improved realistic modelling leads to ~ I order of magnitude loss of sensitivity (unfortunately).

- Using this realistic modelling: optimization of the experiment: work in progress
- New runs in the 10-20 GHz frequency range performed at CEA
- Improvements:
  - new low-noise amplifiers
  - new spectrum analyses

### Scalar Ultra Light Dark Matter: the dilaton

Remark: the QCD coupling of the axion implies that the pions mass depends quadratically on the axion.

see Kim and Perez, arXiv 2023

 $\Rightarrow$  Part of the following discussion can be extended to the axion

# A scalar DM is expected to break the equivalence principle

• An effective Lagrangian for the scalar-matter coupling

$$\mathcal{L}_{\text{mat}}\left[g_{\mu\nu},\Psi,\varphi\right] = \mathcal{L}_{SM}\left[g_{\mu\nu},\Psi\right] + \varphi^{i} \left[\frac{d_{e}^{(i)}}{4e^{2}}F_{\mu\nu}F^{\mu\nu} - \frac{d_{g}^{(i)}\beta_{3}}{2g_{3}}F_{\mu\nu}^{A}F_{A}^{\mu\nu} - \sum_{j=e,u,d}\left(\frac{d_{m_{j}}^{(i)}}{4m_{j}} + \gamma_{m_{j}}d_{g}^{(i)}\right)m_{j}\bar{\psi}_{j}\psi_{j}\right]$$

- Couplings usually considered:
  - linear in  $\varphi$ : lowest order expansion (cfr Damour-Donoghue)
  - quadratic in  $\varphi$ : lowest order if there is a Z<sub>2</sub> symmetry (cfr Stadnik et al)
- This leads to a space-time dependance of some constants of Nature to the scalar field  $lpha(arphi) = lpha \left(1 + d_e^{(i)} \varphi^i\right)$

$$m_{j}(\varphi) = m_{j} \left( 1 + \frac{d_{m_{j}}^{(i)} \varphi^{i}}{M_{3}(\varphi)} \right) \quad \text{for } j = e, u, d$$
$$\Lambda_{3}(\varphi) = \Lambda_{3} \left( 1 + \frac{d_{g}^{(i)} \varphi^{i}}{g} \right)$$

#### Can be interpreted as a signature of a violation of the Einstein Equivalence Principle: oscillations of the constants of Nature!

see also Arvanitaki et al, PRD 2015, Hees et al, PRD, 2018

see Damour and Donoghue, PRD, 2010

#### Two experiments developed at SYRTE

# Search for a periodic signal in Cs/Rb comparison

 Cs/Rb FO2 atomic fountain data from SYRTE: high accuracy and high stability, data used from 2008

see J. Guéna et al, Metrologia, 2012 and J. Guéna et al., IEEE UFFC, 2012

• Search for a periodic signal in the data using Scargle's method, see Scargle ApJ, 1982





A. Hees, J. Guéna, M. Abgrall, S. Bize, P. Wolf, PRL, 2016

### Search for a periodic signal in a Mach-Zender interferometer

• New type of experiment proposed. Simplified principle:



- Interpretation: comparison of an atomic frequency with itself in the past
- Main advantage: explored frequency range ~ kHz-MHz while standard clocks are limited to 100 mHz



Systèmes de Référence Temps-Espace

### The DAMNED experiment (DArk Matter from Non Equal Delays)



- the "clock" is a laser cavity (both length and laser frequency oscillate)
- the length of the fiber oscillates
- the refractive index of the fiber oscillates
- First experiment built @SYRTE (E. Savalle's PhD with P-E Pottie, F. Franck, E. Cantin) and data analyzed taken into account the stochasticity of the signal
  - no significant periodic signal is detected in the 10-200 kHz frequency band

Observatoire SYRTE

Systèmes de Référence Temps-Espac

#### Constraints on the linear couplings

Assuming the DM density to be constant over the whole Solar System (0.4 GeV/cm<sup>3</sup>)



#### Atom interferometers are sensitive to such DM candidates as well

- work from J. Gué, PhD student
- Calculations performed following method from Storey and Cohen-Tannoudji, J. Phys, 1994. Exemple for a Mach-Zender:



- Dilaton DM field impacts:
  - Classical trajectories of atoms
  - Rest mass/transition energy (Lagrangian + kick velocity)
  - Laser reference and frequency

### Phase shift induced by DM in various AI setup and sensitivity of various experiments

work from J. Gué, PhD student

 Standard Mach-Zender: used in Standford with <sup>85</sup>Rb and <sup>87</sup>Rb and for a gravimeter in Wuhan using <sup>87</sup>Rb

see P.Asenbaum et al, PRL, 2020 for standford and Z. Hu et al, PRA, 2020 for Wuhan

 Future AION10 gradiometer: 2 Mach-Zender with Large Momentum Transfer stacked at different elevations

see e.g. Badurina et al, PRD, 2022

 Future MAGIS-like experiment: 2 colocated Mach-Zender with Large Momentum Transfer using 2 isotopes: advantageous for **UFF** tests

see e.g. Abe et al, Quantum Sc. and Tech., 2021

#### Results: search for dilaton



#### Results: search for axion work from J. Gué, PhD student

• The mass of the pion oscillate due to the QCD coupling of



Low sensitivity from Stanford and Wuhan, good sensitivity of AION10 and even better with MAGIS-like scenario

#### Conclusion

- Searches for Dark Matter of mass < I eV (bosonic) is very active
- Several models exist: scalar field, axion, dark photon, ... with different phenomenology: oscillations (possible screening), topological default, ...
- We (SYRTE + collaborations) are involved in:
  - theoretical exploration, predictions of such models
  - proposition of new experiments
  - accurate modelling, optimization of existing experiments
  - perform some experiments
  - dedicated data analysis (sometimes tricky: stochasticity of signal)
- Very recent results: new proposal for an experiment to search for DP
  - modelling/optimisation of SHUKET (DP)
  - impact of dilatons/axions on UFF measurements and AI
  - GASTON: search for transient DM candidate with Galileo

### Are there other signatures to be searched in lab data that can help constraining DE models?



## The field has a frequency distribution due to the DM velocity distribution

• The oscillation frequency depends on the velocity



See Centers et al, arXiv1905.13650 and Foster et al, PRD, 2018 Savalle et al, PRL 2021 Linear and quadratic couplings have a different phenomenology

• Linear coupling

$$\varphi^{(1)}(t, \boldsymbol{x}) = \varphi_0 \cos\left(\boldsymbol{k} \cdot \boldsymbol{x} - \omega t + \delta\right) - s_A^{(1)} \frac{GM_A}{c^2 r} e^{-r/\lambda_{\varphi}}$$

DM, atomic sensors are more sensitive A fifth force generated by a body - UFF tests are more sensitive

Quadratic coupling: no more Yukawa interaction, richer phenomenology

Can be screened or enhanced (scolarisation)

Both atomic sensors and UFF tests are sensitive to this behaviour

Linear and quadratic couplings have a different phenomenology

• Linear coupling

$$\varphi^{(1)}(t, \boldsymbol{x}) = \varphi_0 \cos\left(\boldsymbol{k} \cdot \boldsymbol{x} - \omega t + \delta\right) - \left(s_A^{(1)} \frac{GM_A}{c^2 r} e^{-r/\lambda_{\varphi}}\right)$$

DM, atomic sensors are more sensitive A fifth force generated by a body - UFF tests are more sensitive

Quadratic coupling: no more Yukawa interaction, richer phenomenology

$$\varphi = \tilde{\varphi}(r)\varphi_0 \cos mt$$

Can be screened or enhanced (scolarisation)

Both atomic sensors and UFF tests are sensitive to this behaviour





Screening for positive couplings and scalarization for negative couplings!

### This leads to a rich phenomenology

Comparison of atomic frequencies:

$$Y(t, \boldsymbol{x}) = K + \left(\Delta \kappa^{(2)} \frac{\varphi_0^2}{2} \left(1 - s_A^{(2)} \frac{GM_A}{c^2 r}\right)^2 + \left(\Delta \kappa^{(2)} \frac{\varphi_0^2}{2} \cos\left(2\omega t + 2\delta\right) \left(1 - s_A^{(2)} \frac{GM_A}{c^2 r}\right)^2 + \left(\Delta \kappa^{(2)} \frac{\varphi_0^2}{2} \cos\left(2\omega t + 2\delta\right) \left(1 - s_A^{(2)} \frac{GM_A}{c^2 r}\right)^2\right)^2 + \left(\Delta \kappa^{(2)} \frac{\varphi_0^2}{2} \cos\left(2\omega t + 2\delta\right) \left(1 - s_A^{(2)} \frac{GM_A}{c^2 r}\right)^2\right)^2 + \left(\Delta \kappa^{(2)} \frac{\varphi_0^2}{2} \cos\left(2\omega t + 2\delta\right) \left(1 - s_A^{(2)} \frac{GM_A}{c^2 r}\right)^2\right)^2 + \left(\Delta \kappa^{(2)} \frac{\varphi_0^2}{2} \cos\left(2\omega t + 2\delta\right) \left(1 - s_A^{(2)} \frac{GM_A}{c^2 r}\right)^2\right)^2 + \left(\Delta \kappa^{(2)} \frac{\varphi_0^2}{2} \cos\left(2\omega t + 2\delta\right) \left(1 - s_A^{(2)} \frac{GM_A}{c^2 r}\right)^2\right)^2 + \left(\Delta \kappa^{(2)} \frac{\varphi_0^2}{2} \cos\left(2\omega t + 2\delta\right) \left(1 - s_A^{(2)} \frac{GM_A}{c^2 r}\right)^2\right)^2 + \left(\Delta \kappa^{(2)} \frac{\varphi_0^2}{2} \cos\left(2\omega t + 2\delta\right) \left(1 - s_A^{(2)} \frac{GM_A}{c^2 r}\right)^2\right)^2 + \left(\Delta \kappa^{(2)} \frac{\varphi_0^2}{2} \cos\left(2\omega t + 2\delta\right) \left(1 - s_A^{(2)} \frac{GM_A}{c^2 r}\right)^2\right)^2 + \left(\Delta \kappa^{(2)} \frac{\varphi_0^2}{2} \cos\left(2\omega t + 2\delta\right) \left(1 - s_A^{(2)} \frac{GM_A}{c^2 r}\right)^2\right)^2 + \left(\Delta \kappa^{(2)} \frac{\varphi_0^2}{2} \cos\left(2\omega t + 2\delta\right) \left(1 - s_A^{(2)} \frac{GM_A}{c^2 r}\right)^2\right)^2 + \left(\Delta \kappa^{(2)} \frac{\varphi_0^2}{2} \cos\left(2\omega t + 2\delta\right) \left(1 - s_A^{(2)} \frac{GM_A}{c^2 r}\right)^2\right)^2 + \left(\Delta \kappa^{(2)} \frac{\varphi_0^2}{2} \cos\left(2\omega t + 2\delta\right) \left(1 - s_A^{(2)} \frac{GM_A}{c^2 r}\right)^2\right)^2 + \left(\Delta \kappa^{(2)} \frac{\varphi_0^2}{2} \cos\left(2\omega t + 2\delta\right) \left(1 - s_A^{(2)} \frac{GM_A}{c^2 r}\right)^2\right)^2 + \left(\Delta \kappa^{(2)} \frac{\varphi_0^2}{2} \cos\left(2\omega t + 2\delta\right) \left(1 - s_A^{(2)} \frac{GM_A}{c^2 r}\right)^2\right)^2 + \left(\Delta \kappa^{(2)} \frac{\varphi_0^2}{2} \cos\left(2\omega t + 2\delta\right) \left(1 - s_A^{(2)} \frac{GM_A}{c^2 r}\right)^2\right)^2$$

Position dependent: clocks on elliptic orbit? Comparison clock in space versus clock on ground?

oscillation, amplitude depends on position

• UFF measurements

$$[\Delta \boldsymbol{a}]_{A-B} = \Delta \bar{\alpha}^{(2)} \frac{\varphi_0^2}{2} \left( 1 - s_C^{(2)} \frac{GM_c}{c^2 r} \right) \left( -\frac{GM_c}{r^3} \boldsymbol{x} \, s_C^{(2)} \right) - \left( \frac{GM_c}{r^3} \boldsymbol{x} \, s_C^{(2)} \right) \cos \left( 2\omega t + 2\delta \right) \\ + \left( \left( 1 - s_C^{(2)} \frac{GM_c}{c^2 r} \right) \omega \boldsymbol{v} \sin \left( 2\omega t + 2\delta \right) \right) \left( -\frac{GM_c}{r^3} \boldsymbol{x} \, s_C^{(2)} \right) + \left( \frac{GM_c}{r^3} \boldsymbol{x} \, s_C^{(2)} \right) \cos \left( 2\omega t + 2\delta \right) \\ + \left( \frac{GM_c}{r^3} \boldsymbol{x} \, s_C^{(2)} \right) \left( -\frac{GM_c}{r^3} \boldsymbol{x} \, s_C^{(2)} \right) \left( -\frac{GM_c}{r^3} \boldsymbol{x} \, s_C^{(2)} \right) \left( -\frac{GM_c}{r^3} \boldsymbol{x} \, s_C^{(2)} \right) \right) \left( -\frac{GM_c}{r^3} \boldsymbol{x} \, s_C^{(2)} \right) \left($$

 $\eta$  that depends on r (directly related to Eöt-Wash and MICROSCOPE results)

2 terms that oscillate, amplitude depends on position

See A. Hees et al, PRD, 2018

They are all sensitive to screening/scalarization

#### Constraints on the quadratic couplings



Impact of screening

Impact of scalarization

Being in space is favorable ! Scalar field tends to vanish at the Earth surface

# A vector DM will interact with electromagnetism

• An effective Lagrangian for the vector-matter coupling

$$\mathcal{L}_{\text{mat}}\left[\Psi, g_{\mu\nu}, X_{\mu}\right] = \mathcal{L}_{\text{SM}}\left[\Psi, g_{\mu\nu}\right] - \frac{\chi}{2} F^{\mu\nu} X_{\mu\nu} + \dots$$

see Horns et al, JCAP, 2013 and references therein

- Kinetic mixing coupling  $\chi$  characterises the coupling with EM
- Other couplings with matter can be considered like to the B-L current: leads to a violation of the UFF
- The hidden photon  $X^{\mu}$  will mix with the usual photon  $A^{\mu}$

$$\Box A^{\mu} = -\chi \Box X^{\mu}$$
$$\Box X^{\mu} + m^2 X^{\mu} = -\chi \Box A^{\mu}$$

A hidden photon field will generate a small EM field and vice versa

# An oscillating DM vector field will generate a small electric field

• Oscillating DM vector field  $\vec{X} = \vec{X}_0 \cos mt$  will generate an EM field

$$\vec{A} = -\chi \vec{X}$$

and in particular a small electric field

$$\vec{E}_{\rm DM} = -\partial_t \vec{A} = -m\chi \vec{X}_0 \sin mt$$

• As a reminder: the amplitude of oscillation is related to the DM energy density  $a \mid \vec{z} \mid^2$ 

$$\rho = \frac{m^2 \left| \vec{X}_0 \right|^2}{2}$$

#### In a DM vector field, a dish antenna will generate an EM field that will be focused in its center

- the electric field // to a conductor surface vanishes (boundary condition)
- The surface of the dish will generate a propagating electric field to vanish the DM electric field
- For a spherical dish, the electric field will be focused at the center + nonrelevant electric field will be focused at the focal point
- Sensitivity

S

center

## A scalar field with a quartic potential can form topological defects



• Width related to the mass of the scalar field  $W \sim 1/m$ 

- Amplitude related to DM energy density  $\varphi_0^2 \sim \rho_{\rm DM}$
- Cross the Earth with a velocity (DM velocity distribution)