

Quantum interference measurement of the free fall of anti-hydrogen

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Introduction - GBAR

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The GBAR Experiment



- Measurement of gravity interaction of matter on anti-matter, test of the equivalence principle
- Measurement of the free fall time of antihydrogen atoms
- Atom-by-atom detection
- Objective of precision on \bar{g} of the order of 10⁻² with 1000 atoms



- Ultra cold \bar{H}^+ in ground state of an ion trap
- Photodetachment of excess positron by laser pulse
- Freefall until annihilation on Micomega plates
- Objective of precision on \bar{g} of the order of 10⁻² with 1000 atoms

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Theoretical Modelisation



- Bounce on the Casimir Polder potential
- We aim for a high reflection probability
- Easier to achieve with fewer bounces
- Reflexion due to high variation of the potentiel near the surface

In the GBAR experiment we hope to drop the \bar{H} from $h \propto$ 10 μ m

[3] G. Dufour et al. Quantum reflection of antihydrogen from the Casimir potential above matter slabs PRA, 2013 3

Introduction - Quantum GBAR





- New measurement method, based on interferences
- Quantum bounce on the attractive Casimir-Polder potential

- Atom-by-atom detection
- Objective of precision on \bar{g}/g of the order of 10⁻⁵ with 1000 atoms^[4]

[4] P-P.Crépin et al. Quantum interference test of the equivalence principle on antihydrogen PRA, 2019

We want to solve the eigen-value equation :

$$-\frac{\hbar^2}{2m}\frac{d^2\psi(z)}{dz^2} + V(z)\psi(z) = E\psi(z) \quad \text{with}: \quad V(z) = \begin{cases} mgz & \text{is } z > 0\\ +\infty & \text{else} \end{cases}$$

 $\simeq 145 \text{Hz}$



• Solved by the Airy function : $\psi_n(z) = \frac{\theta(z)}{\sqrt{l_g} \text{Ai}'(-\lambda_n)} \text{Ai}\left(\frac{z}{l_g} - \lambda_n\right)$ • Where $l_g := \left(\frac{\hbar^2}{2m^2 g}\right)^{\frac{1}{3}} \simeq 5.871 \,\mu\text{m } \&$ λ_n are the zeros of the Airy function • Energy scale : $\epsilon_q = m g \, l_g \simeq 0.6 \text{ peV}$

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Introduction - Numerical Simulation & Measurement $\sqrt{1 - KB}$



- Random draw of N point in the simulated current with $g_0 = 9.81$
- Simulation of the current for different values of *g*
- Maximum Likelihood estimator ; \hat{g} ; over all the simulations
- Standard deviation of \hat{g} after M repetitions gives the measurment

We have $\bar{g}/g \propto 1e-5$



Toward fewer bounce

Result - Two waves interference

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- The fewer the bounces, the simpler the interference pattern
- Possible to test the interferometer with a Hydrogen atom beam
- A two wave interference regime, between one bounce and zero bounce



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Result - Focalisation of the wavepacket

Study of the one bounce regime



• After only one bounce the wavepacket refocalise

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- This focal point is not well defined in space
- We can use this to create interferences inside the bounce
- Considering the miror as a lens we can determine the focal time as $t_f = \frac{t_i^2}{t_i - \frac{v_i}{2g}}$ Where t_i and v_i are the impact time and velocity



The quantum evolution shows explicitly the interferences created at the focal spot

1/LKB



The quantum evolution shows explicitly the interferences created at the focal spot

1/LKB



The quantum evolution shows explicitly the interferences created at the focal spot

1 LKB



- $\cdot\,$ Cut at the focal time
- Simple interference pattern
- Maximal contrast
- Can be predicted by the stationnary phase method over the action

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Result - Numerical estimation of the precison



Take Home Message : With about 100 atoms we can achieve a relative precision of the 10⁻⁵ order and have a much simpler interference pattern while having shorten the experimentation time.

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- We proposed a new kind of atomic interferometer
- It can be applied beyond the scope of the GBAR experiment
- The GRASIAN experiment aims to test the quantum bounce of Hydrogen
- We hope to apply this method to exotic atoms with a very short lifespan atoms or a very little sample
- We are currently working on the full model with loss and photodetachment taken into account